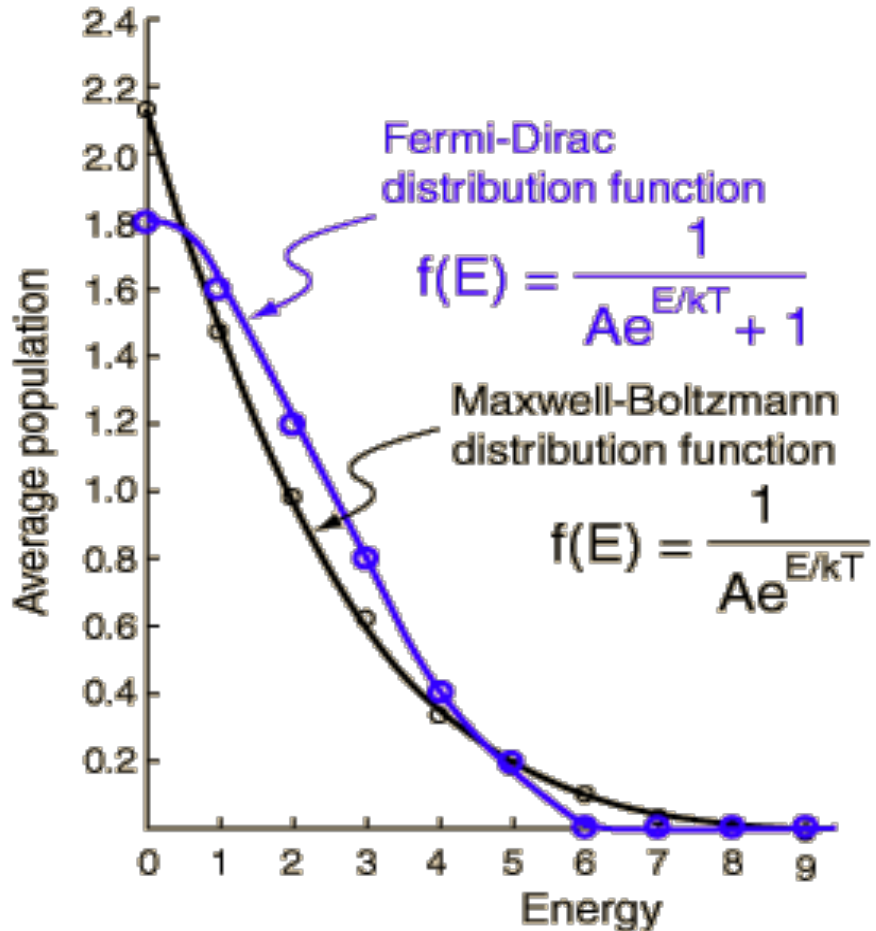
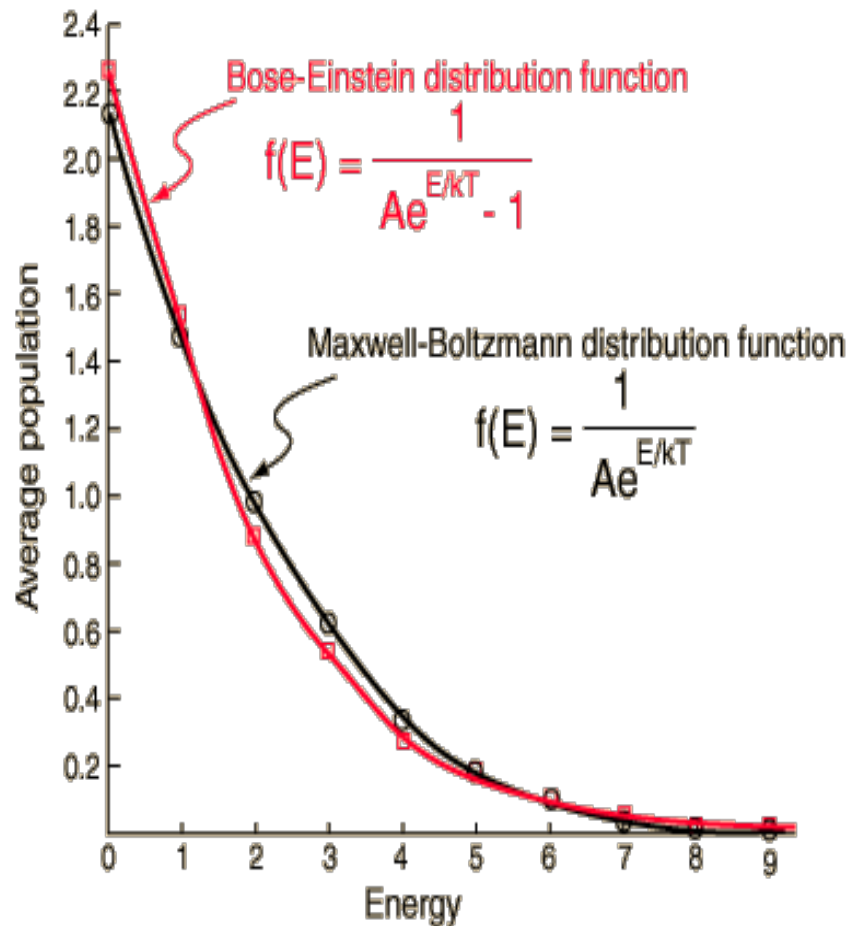


Modern Physics (Phys. IV): 2704

Professor Jasper Halekas
Van Allen 70
MWF 12:30-1:20 Lecture

Classical/Quantum Energy Distribution Functions



Degeneracy/Density of States

- d_n is the degeneracy (or maximum occupancy) of the energy level E_n
 - E.g. in the hydrogen atom the electron subshell $n = 2, l = 1$ has degeneracy $d_n = 6$
- $g(E)$ is the density of states per unit volume
 - $V g(E)$ is the continuous equivalent of d_n

Discrete Energy Distribution

$$N_n = d_n P_n$$

d_n = degeneracy of level n

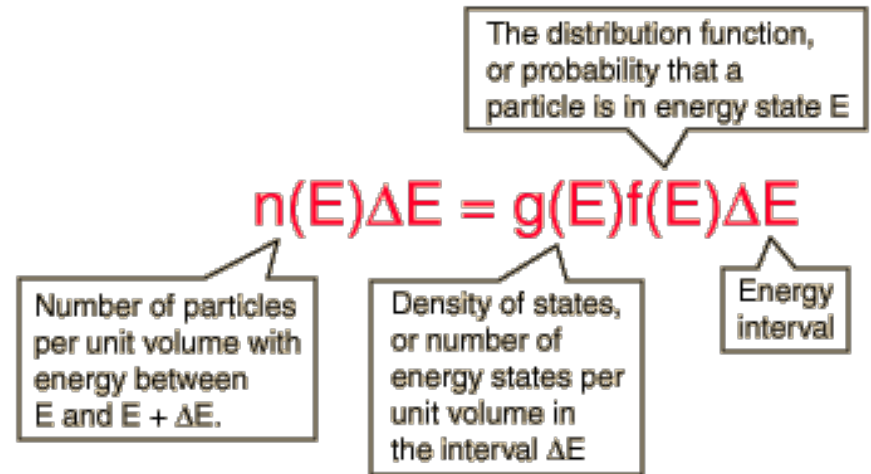
P_n = probability of energy n

$$N = \sum_n N_n = \sum_n d_n P_n$$

$$\langle E \rangle = \frac{\sum_n N_n E_n}{N} = \frac{\sum_n d_n P_n E_n}{\sum_n d_n P_n}$$

Populated States

The number of populated states per unit volume $n(E)$ is proportional to the product of the density of states $g(E)$ and the energy distribution function $f(E)$



Concept Check

- What is the correct normalization condition for N particles?
- $\text{Integral}(V g(E) f(E) dE) = 1$
- $\text{Integral}(g(E) f(E) dE) = 1$
- $\text{Integral}(V g(E) f(E) dE) = N$
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Continuous Energy Distribution

$$dN = N(E) dE$$

$$= V g(E) f(E) dE$$

V = volume

$g(E)$ = density of states

⊗ Energy E per volume

$f(E)$ = energy distribution function
= probability of energy E

$$N = \int dN = \int_0^{\infty} N(E) dE$$

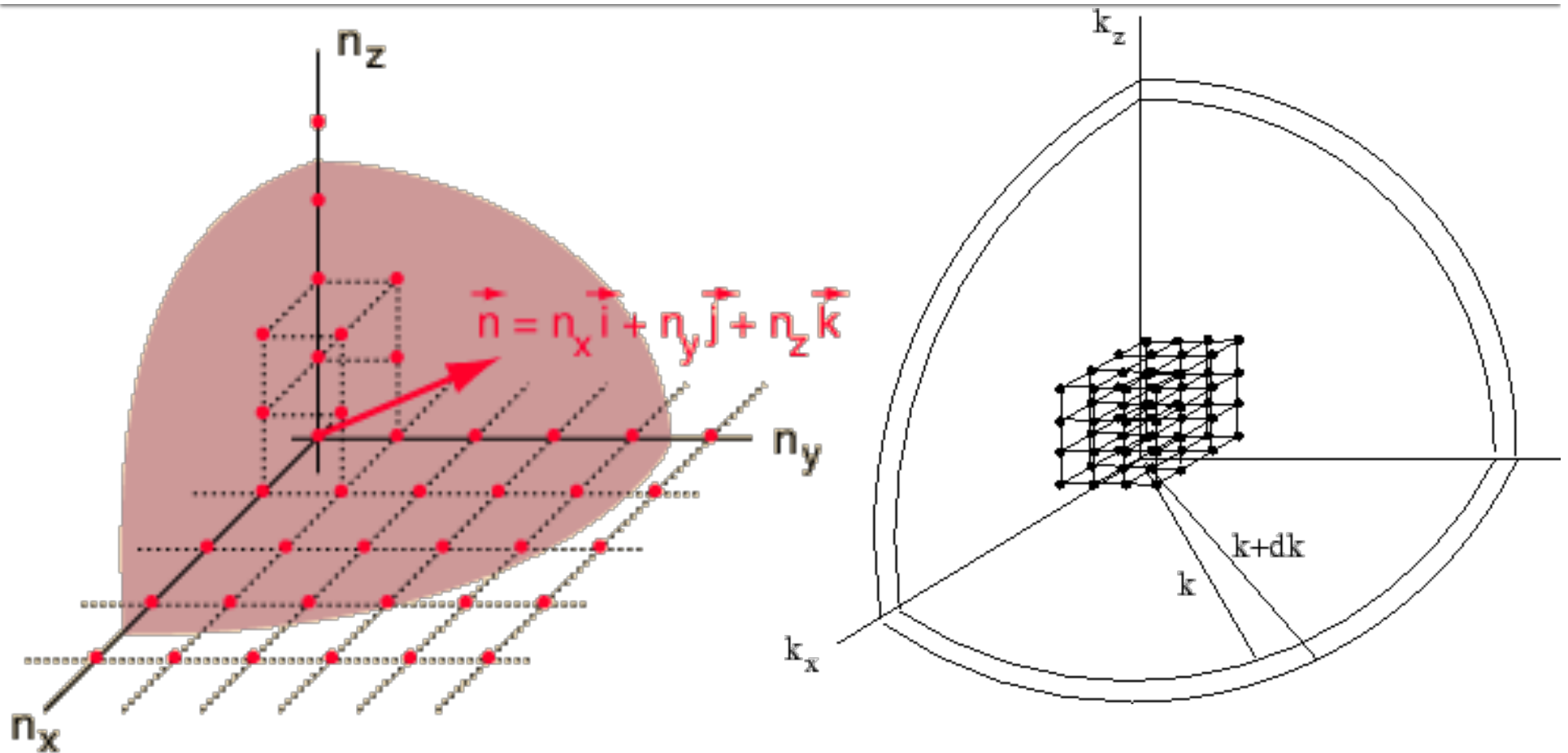
$$= \int_0^{\infty} V g(E) f(E) dE$$

$$\langle E \rangle = \frac{\int_0^{\infty} E N(E) dE}{N}$$

$$= \frac{V \int_0^{\infty} E g(E) f(E) dE}{V \int_0^{\infty} g(E) f(E) dE}$$

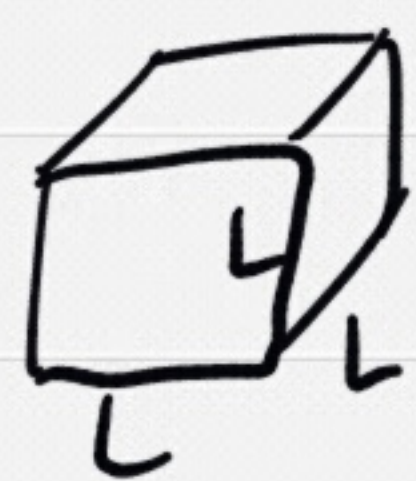
$$= \frac{\int_0^{\infty} E g(E) f(E) dE}{\int_0^{\infty} g(E) f(E) dE}$$

Density of States in a Gas in a Box



Density of States: Gas

Wave function for particle in box



$$\Psi(x, y, z) = A \sin\left(\frac{n_x \pi x}{L}\right) \times \sin\left(\frac{n_y \pi y}{L}\right) \times \sin\left(\frac{n_z \pi z}{L}\right)$$

$$E = \frac{\hbar^2 \pi^2}{2mL^2} (n_x^2 + n_y^2 + n_z^2)$$
$$= \frac{\hbar^2}{2m} (k_x^2 + k_y^2 + k_z^2)$$

- # of states from E to $E + dE$
- A shell in n_x, n_y, n_z space w/ all positive

$g(n) dn$ proportional to

$$\frac{1}{8} 4\pi n^2 dn$$

\uparrow \uparrow \nwarrow
 $n_x, n_y, n_z > 0$ Area of shell thickness of shell

- spin s particles $\Rightarrow 2s+1$ orientations

$$g(n) dn = \frac{1}{8} \frac{2s+1}{V} 4\pi n^2 dn$$

$$g(E) dE = g(n) dn$$

$$g(E) = g(n) / \left(\frac{dE}{dn} \right)$$

$$E = \frac{\hbar^2 \pi^2 n^2}{2mL^2} \quad \Leftrightarrow \quad n = \sqrt{\frac{2mL^2 E}{\hbar^2 \pi^2}}$$

$$\frac{dE}{dn} = \frac{\hbar^2 \pi^2 n}{mL^2}$$

$$\Rightarrow g(E) = \frac{1}{8} \frac{2s+1}{V} 4\pi n^2 / \left(\frac{\hbar^2 \pi^2}{mL^2} n \right)$$

$$= \frac{1}{8} \frac{2s+1}{V} 4\pi n / \left(\frac{\hbar^2 \pi^2}{mL^2} \right)$$

$$= \frac{1}{8} \frac{2s+1}{V} \cdot 4\pi \cdot \sqrt{\frac{2mL^2 E}{\hbar^2 \pi^2}} / \left(\frac{\hbar^2 \pi^2}{mL^2} \right)$$

$$= \frac{1}{8} \frac{2s+1}{V} \cdot 4\pi \frac{\sqrt{2m} L \sqrt{E}}{\hbar^2 \pi} \cdot \frac{mL^2}{\hbar^2 \pi^2}$$

$$= \left(\frac{2s+1}{\sqrt{2}} \frac{m^{3/2}}{\pi^2 \hbar^3} \sqrt{E} \right)$$

Note L^3 cancels \checkmark

Density of States: Photons

Same derivation, but

$$2s+1 \rightarrow 2 \text{ (RH or LH)}$$

$$E = h\nu = \hbar\omega$$
$$= \hbar kc$$

$$= \frac{\hbar c \pi n}{L} \Leftrightarrow n = \frac{E L}{\hbar \pi c}$$

$$dE/dn = \hbar c \pi / L$$

$$g(n) = 2 \cdot \frac{1}{8} \cdot 4\pi n^2 \cdot \frac{1}{V}$$

$$g(E) = g(n) / dE/dn$$

$$= \frac{\pi n^2}{V} \cdot \frac{L}{\hbar c \pi}$$

$$= \frac{\pi}{V} \frac{E^2 L^2}{\hbar^2 \pi^2 c^2} \cdot \frac{L}{\hbar c \pi}$$

$$= \frac{E^2}{\pi^2 \hbar^3 c^3}$$

- Fills hole in our derivation of Planck blackbody formula

Maxwell - Boltzmann Distribution

$$f(E) = \frac{1}{A} e^{-E/kT}$$

$$g(E) = \frac{2s+1}{\sqrt{2}} \frac{m^{3/2}}{\pi^2 \hbar^3} \sqrt{E}$$

$$N(E) = V g(E) f(E)$$

$$= \frac{1}{A} \cdot \frac{2s+1}{\sqrt{2}} \frac{m^{3/2}}{\pi^2 \hbar^3} \cdot V \cdot \sqrt{E} e^{-E/kT}$$

$$\int_0^{\infty} N(E) dE = N$$

$$\int_0^{\infty} \sqrt{E} e^{-E/kT} dE = \frac{\sqrt{\pi}}{2} (kT)^{3/2}$$

$$\Rightarrow \frac{1}{A} \frac{2s+1}{\sqrt{2}} \frac{m^{3/2}}{\pi^2 \hbar^3} \cdot V \cdot \frac{\sqrt{\pi}}{2} (kT)^{3/2} = N$$

$$\Rightarrow N(E) = \frac{2N}{\sqrt{\pi} (kT)^{3/2}} \sqrt{E} e^{-E/kT}$$

M-B speed distribution

$$N(v) = N(E) dE/dv = N(E) \cdot mv$$

$$= \frac{2N}{\sqrt{\pi} (kT)^{3/2}} \cdot \sqrt{\frac{1}{2} mv^2} \cdot mv \cdot e^{-\frac{1}{2} mv^2/kT}$$

$$= N \sqrt{\frac{2}{\pi}} \left(\frac{m}{kT}\right)^{3/2} \cdot v^2 e^{-\frac{mv^2}{2kT}}$$

M-B velocity distribution

$$\begin{aligned}\int_0^\infty N(v) dv &= \iiint N(v_x, v_y, v_z) dv_x dv_y dv_z \\ &= \iiint N(v, \theta, \varphi) \cdot v^2 dv \sin\theta d\theta d\varphi \\ &= 4\pi \int_0^\infty N(v, \theta, \varphi) \cdot v^2 dv\end{aligned}$$

$$\Rightarrow N(v, \theta, \varphi) = N(v_x, v_y, v_z)$$

$$= N(v) \cdot \frac{1}{4\pi v^2}$$

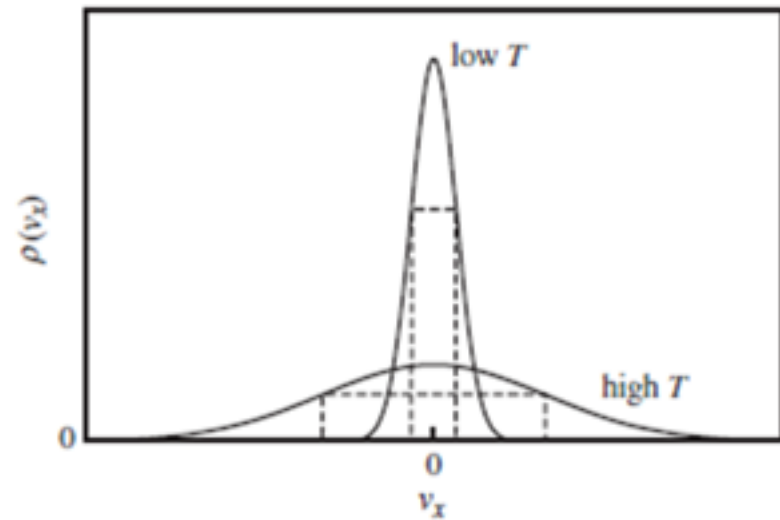
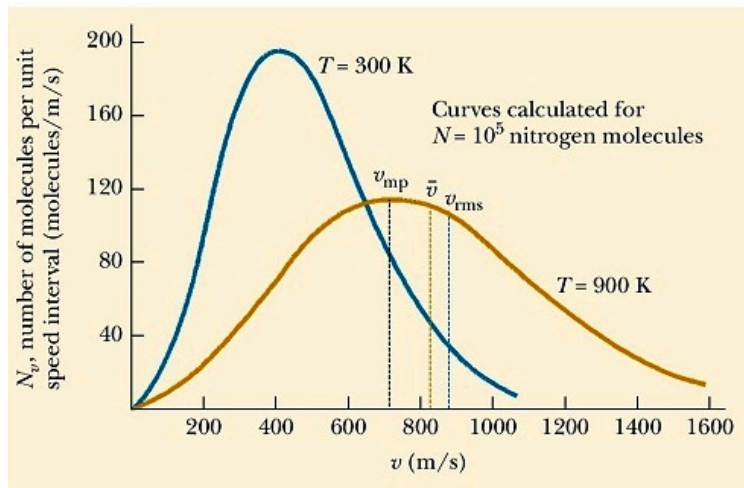
$$= N \cdot \sqrt{\frac{2}{\pi}} \cdot \frac{1}{4\pi} \cdot \left(\frac{m}{kT}\right)^{3/2} \cdot v^2 e^{-mv^2/2kT} \cdot \frac{1}{v^2}$$

$$= \boxed{N \left(\frac{m}{2\pi kT}\right)^{3/2} e^{-m(v_x^2 + v_y^2 + v_z^2)/2kT}}$$

1-d equivalent

$$N(v_x) = N \sqrt{\frac{m}{2\pi kT}} e^{-mv_x^2/2kT}$$

Maxwell-Boltzmann Distribution



Doppler Broadening

put $v_x =$ "line-of-sight" velocity

Doppler shift

$$f = \sqrt{\frac{1 - v_x/c}{1 + v_x/c}} f_0$$

$$\sim f_0 (1 - v_x/c) \quad |v_x| \ll c$$

$$\Rightarrow v_x = c (1 - f/f_0)$$

$$N(f) = N(v_x) / |df/dv_x|$$

$$= N(v_x) / (f_0/c)$$

$$= N(v_x) \cdot c/f_0$$

$$= \frac{Nc}{f_0} \sqrt{\frac{m}{2\pi kT}} e^{-mc^2(1-f/f_0)^2/2kT}$$

$$\text{FWHM } \Delta f = 2f_0 \sqrt{(2 \ln 2) kT/mc^2}$$

Doppler Broadening

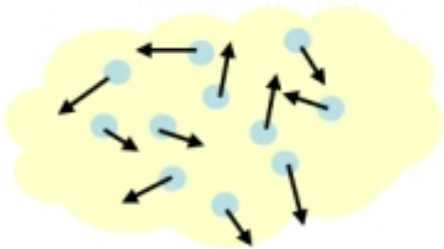
Gas particles at rest



Emission line spectrum with narrow lines



Gas particles with random motions



Emission line spectrum with thermal line broadening

