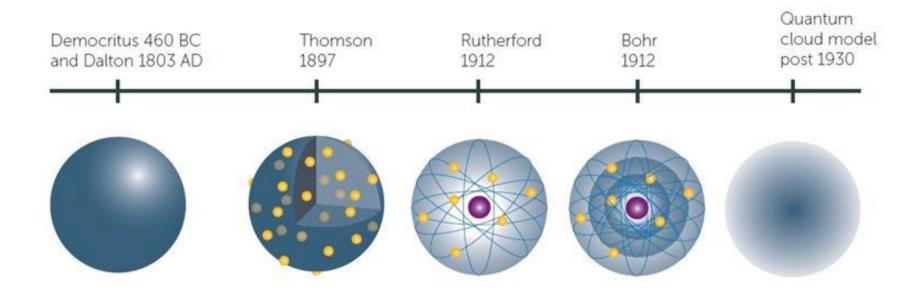


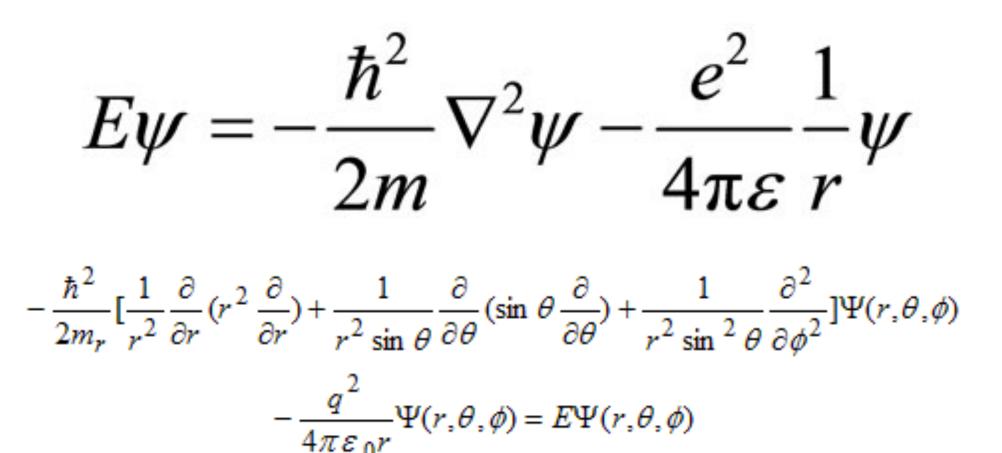
## Modern Physics (Phys. IV): 2704

Professor Jasper Halekas Van Allen 70 MWF 12:30-1:20 Lecture

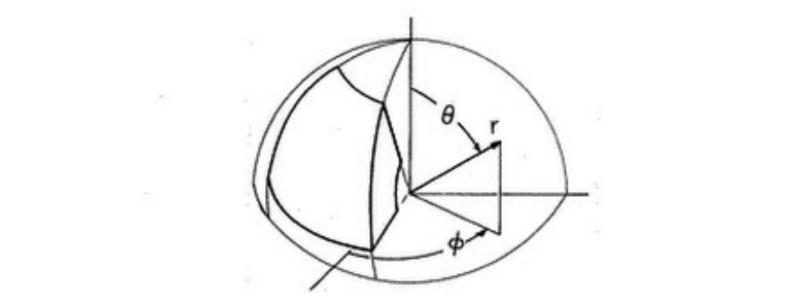
#### **Models of the Atom**



#### Schrödinger Equation for Hydrogen



#### Laplacian in Spherical Coordinates



$$\nabla^2 = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial}{\partial r} \right) + \frac{1}{r^2 \sin(\theta)} \frac{\partial}{\partial \theta} \left( \sin(\theta) \frac{\partial}{\partial \theta} \right) + \frac{1}{r^2 \sin^2(\theta)} \frac{\partial^2}{\partial \varphi^2}$$

#### **Concept Check**

- Given an equation f(x) = g(y) that is true for all values of the independent variables x and y, what \*must\* be true of the functions f and g?
- A. Both functions are equal to zero
- B. Both functions are polynomials
- C. Both functions are constant
- D. This equation cannot be solved

#### **Concept Check**

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Hydrogen Atom Time - Independent Schrödinger Eq.  $\begin{aligned} & |-d: -\frac{\hbar^{2}}{2m} \frac{d^{2}\psi}{dx^{2}} + U(x)\Psi(x) = E\Psi(x) \\ & 3-d: -\frac{\hbar^{2}}{2m} P^{2}\psi + U(F)\Psi(F) = E\Psi(F) \end{aligned}$ Cartesian:  $\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}$ spherical:  $\nabla^2 = \frac{1}{\sqrt{2}} \frac{\partial x}{\partial x} (\frac{v^2}{\sqrt{2}})$   $+ \frac{1}{\sqrt{r}\sin\theta} \frac{\partial x}{\partial \phi} (\frac{x^2}{\sqrt{2}})$   $+ \frac{1}{\sqrt{r}\sin\theta} \frac{\partial^2}{\partial \phi^2}$ Hatom  $u(r) = -\frac{e}{4\pi_0}r$  $= -\frac{\hbar^{2}}{(2mr^{2})} \left[ \frac{\sqrt{2}r}{r^{2}} \left( \frac{r^{2}}{r^{2}} + \frac{\sqrt{2}r}{r^{2}} \right) + \frac{\sqrt{2}r}{r^{2}} + \frac{\sqrt{2}r}{r^{2}}$ Use separation of Variables

 $Y(r, q, q) = R(r) \Theta(q) \overline{\Phi}(q)$ 

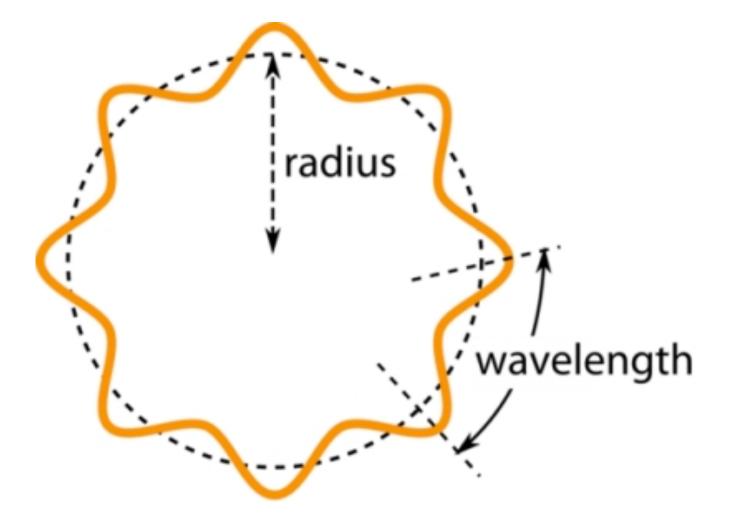
- Divide by RAD - Multiply by 2mr2/22  $-\left[\frac{1}{R}\sqrt{3}r\left(r\lambda\partial R\right)r\right] + \frac{1}{\Theta \sin \theta}\sqrt{3}\sigma(\sin \theta \partial \beta \rho) + \frac{1}{4}\frac{1}{5}\frac{1}{\sin^{2}\theta}\sqrt{3}\rho^{2} - \frac{me^{2}r}{\pi e^{2}r} = \frac{2mE}{4}r^{2}$  $= \frac{1}{2} \left[ \frac{1}{2} \left( r^2 \frac{\partial R_{ar}}{\partial r} \right) + \frac{me^{lr}}{2\pi s_{atl}} - \frac{lmE}{4L} r^2 \right]$  $= - \left[ \frac{\partial s_{ine}}{\partial s_{ine}} \frac{\partial A}{\partial a} \left( s_{ine} \frac{\partial R_{ar}}{\partial a} \right) + \frac{me^{lr}}{4s_{inl}} \frac{\partial L}{\partial a} \right]$ 

- LHS depends only on n - RHS depends on Opp - True for all n Op => both sides must be constant Write constant = Cr  $\Rightarrow \overline{As'nA} \overline{Ao}(s'nA \overline{AA}) + \overline{\pm sin^2 a} \overline{Aa} \overline{Aa} = -Ce$ 

 $= \frac{\sin \phi}{\partial \phi} \left( \sin \phi \partial \phi \right) + \frac{1}{2} \frac{\partial^2 \phi}{\partial \phi^2} = -C_{e} \sin^2 \phi$  $= \frac{\sin \pi}{2} \frac{1}{2} \frac{1}{2}$ LHS depends on of RHS on \$9, so both are constant

- Write second constant = (m - Three OPES for R, Q, I Kdu(ridkir) + mein + 2mE 41 (ridkir) + 2TTEOHI + 41 r2 = Ce Sind de (sine de) + Lesine = Cm  $-/= 2 \frac{1}{2} \frac{1}{2$ Azimuthal Wave Function 1 22 Jos = - (mJ Rewrite (m = mi De Flar = -m2 F => I (p) = A e me To satisfy continuity  $\overline{T}(p) = \overline{T}(p + 2\pi)$  $=)[m_{e} = 0, \pm 1, \pm 2, ---]$ Normalitation: Sul []2dep = So Ae e e dep = 2TA2 = 1 =)  $\overline{T}(\varphi) = \overline{\sqrt{2\pi}} e^{im_{e}\varphi}$ 

#### **Azimuthal Wave Function**

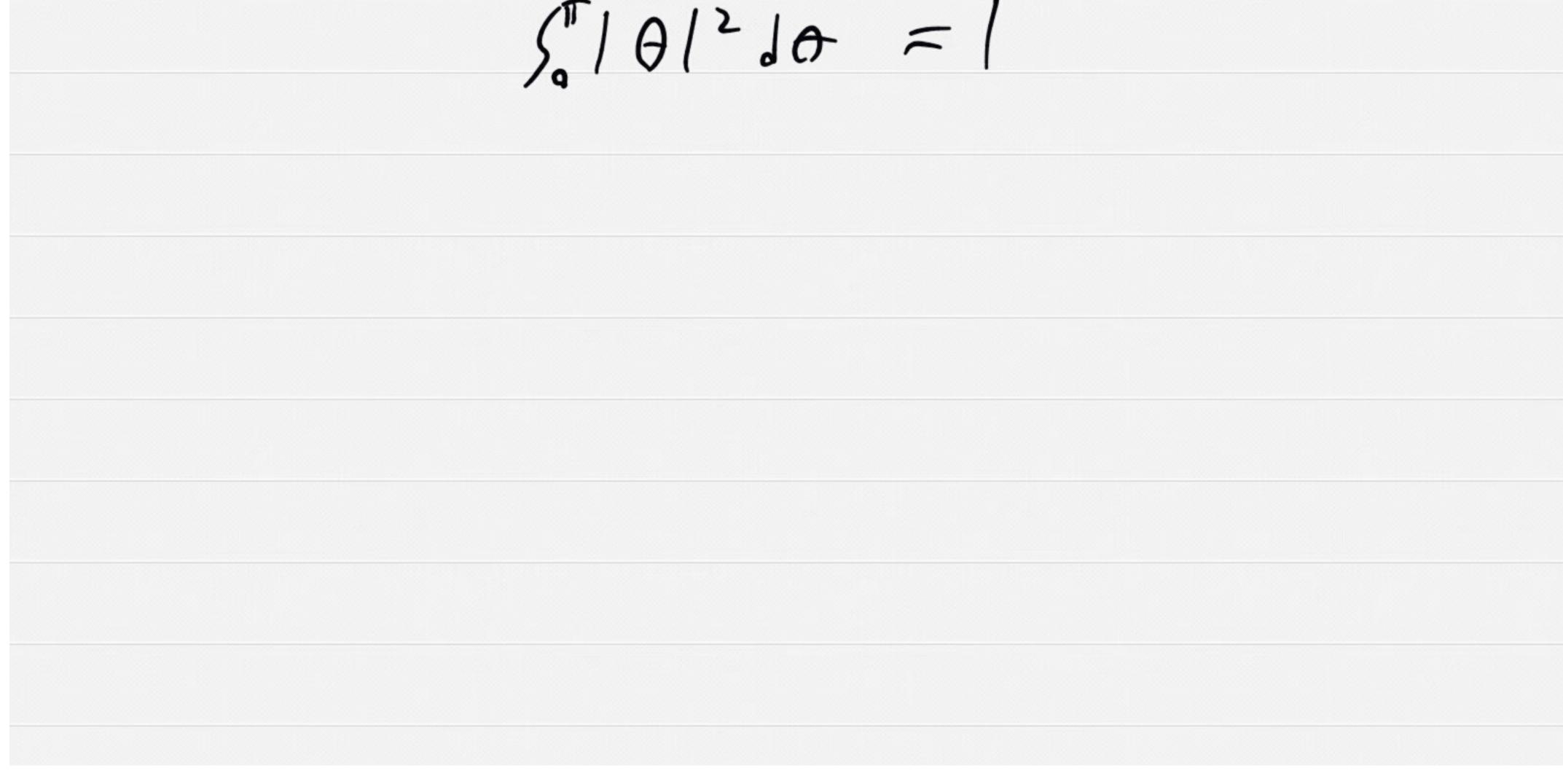


# Azimuthal symmetry of wave functions

- Azimuthal wave function e<sup>imφ</sup>
- Azimuthal probability density  $e^{im\phi}e^{-im\phi} = 1$
- Probability of finding electron is azimuthally symmetric for pure wave function
- Same as for the circular orbits of the Bohr model!

[Polan Wave Function]  $\frac{\sin \theta}{\partial \theta} \frac{\partial \theta}{\partial \theta} \left( \sin \theta \frac{\partial \theta}{\partial \theta} \right) + \left( \frac{2 \sin^2 \theta}{2 \sin^2 \theta} + \frac{2 \sin^2 \theta}{2 \sin^2 \theta} \right) = \frac{2 \sin^2 \theta}{2 \sin^2 \theta}$ =) the do(sind do/da) - me a + Ce A=0 Salutions are APr (cost)

- Per are Legendre Polynomials - For valid solution  $\begin{pmatrix} e = e(e+1) \\ and me = -e - e+1, \dots, e-1, e \end{pmatrix}$  $\Theta(\sigma) = A P_{\mathcal{L}}^{m}(\sigma, \sigma)$ (A determined by normalization



## Legendre Polynomials

P<sup>m</sup>(x) <sub>↑</sub> -0.8-0.6-0.4-0.2 02,04,0608 x -1

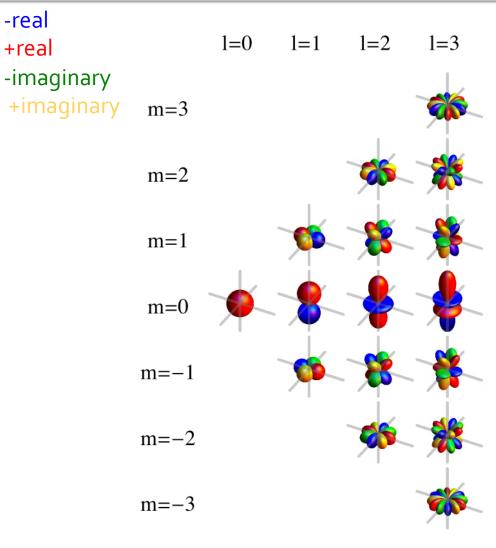
associated legendre functions (normalized)

Table 6-2

The First Few Associated Legendre Functions  $P_{I}^{|m|}(x)$ 

$$\begin{array}{rcl} & P_0^0(x) = 1 \\ = 0 & P_1^0(x) = x = \cos\theta \\ = 1 & \\ = 2 & P_1^1(x) = \sqrt{1 - x^2} = \sin\theta \\ = 3 & P_2^0(x) = 1/2(3x^2 - 1) = 1/2(3\cos^2\theta - 1) \\ = 4 & P_2^1(x) = 3x\sqrt{1 - x^2} = 3\cos\theta \sin\theta \\ & P_2^2(x) = 3(1 - x^2) = 3\sin^2\theta \\ & P_3^0(x) = 1/2(5x^3 - 3x) = 1/2(5\cos^3\theta - 3\cos\theta) \\ m = 0 & P_3^1(x) = 3/2(5x^2 - 1)(1 - x^2)^{1/2} = 3/2(5\cos^2\theta - 1)\sin\theta \\ m = 1 & P_3^2(x) = 15x(1 - x^2) = 15\cos\theta\sin^2\theta \\ m = 3 & P_3^3(x) = 15(1 - x^2)^{3/2} = 15\sin^3\theta \\ m = 4 & \\ \hline \end{array}$$

#### **Spherical Harmonics**



-real

+real

l	$m_{\ell}$	$Y_{\ell m_l}(\theta,\phi) = \Theta_{\ell m_l}(\theta) \Phi_{m_l}(\phi)$
0	0	$(1/4\pi)^{1/2}$
1	0	$(3/4\pi)^{1/2}\cos\theta$
1	$\pm 1$	$\mp (3/8\pi)^{1/2}\sin\theta e^{\pm i\phi}$
2	0	$(5/16\pi)^{1/2}(3\cos^2\theta - 1)$
2	$\pm 1$	$\mp (15 / 8\pi)^{1/2} \sin\theta \cos\theta e^{\pm i\phi}$
2	$\pm 2$	$(15/32\pi)^{1/2}\sin^2\theta e^{\pm 2i\phi}$
2 2	0 ±1	$(5/16\pi)^{1/2}(3\cos^2\theta - 1)$ $\mp (15/8\pi)^{1/2}\sin\theta\cos\theta e^{\pm i\theta}$

$$\Phi_{m\ell}(\phi) = \frac{1}{\sqrt{2\pi}} e^{im_{\ell}\phi}$$
$$\Theta_{\ell m_{\ell}}(\theta) = \left[\frac{2\ell+1}{2} \frac{(\ell-m_{\ell})!}{(\ell+m_{\ell})!}\right]^{1/2} P_{\ell}^{m_{\ell}}(\theta)$$

 $P_{\ell}^{m_{\ell}}(\theta) = associated \ Legendre \ polynomial$ 

### **Angular Probability Density**

